# Oxygen ion conduction in $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> doped with Sb<sub>2</sub>O<sub>3</sub>

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Electrical conduction in b c c  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> doped with Sb<sub>2</sub>O<sub>3</sub> was investigated by measuring electrical conductivity, as a function of temperature and oxygen partial pressure ( $P_{O_2}$ ), and ionic transference number. The  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> doped with 1 to 3 mol % Sb<sub>2</sub>O<sub>3</sub> was stable up to 550° C and showed an oxygen ionic conduction in the  $P_{O_2}$  region of 10<sup>5</sup> to 10<sup>-9</sup> Pa. As the Sb<sub>2</sub>O<sub>3</sub> content increased, ionic conductivity increased up to 2.5 mol % Sb<sub>2</sub>O<sub>3</sub> (1.8 × 10<sup>-3</sup>  $\Omega^{-1}$  cm<sup>-1</sup> at 500° C) and then decreased. However, the activation energy for ionic conduction remained almost unchanged. It was proposed that the  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> contains a lot of oxygen vacancies and incorporated Sb<sup>5+</sup> ions at tetrahedral sites which affect the concentration of oxygen vacancy effective for conduction.

# 1. Introduction

The body-centred cubic  $\gamma$ -form of pure Bi<sub>2</sub>O<sub>3</sub> is a metastable form which appears during the cooling of the high-temperature stable fcc  $\delta$ -form. The bcc  $\gamma$ -form can be stabilized by the addition of small amounts of dopant oxides to form the sillenite group of compounds  $6Bi_2O_3 \cdot MO_x$  in the I23 space group [1-3]. The structure of the sillenite compounds has been determined: the M cations are tetrahedrally coordinated and occupy the body centred sites and the cube corners, linked by hepta-coordinated bismuth atoms. This lattice only achieves the perfect structure where the M cation is tetravalent giving the general formula  $6Bi_2O_3 \cdot MO_2(Bi_{24}M_2O_{40})$  in the unit cell) [4–6].

The pure  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> was first proposed by Sillen [1] to be isomorphous with the bcc  $Bi_{24}Si_2O_{40}$  where bismuth atoms occupy the silicon positions in the cell. Other studies suggested that the  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> contains  $Bi_{26}O_{39}$  in the unit cell [7, 8]. Levin and Roth [2] indicated that the pure  $\gamma\text{-}\text{Bi}_2\text{O}_3$  has the largest cell constant of the series of sillenite compounds reported. Craig and Stephenson [9] suggested that the  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> is isomorphous with  $Bi_{24}^{3+}Bi^{5+}Fe_{1}^{3+}O_{40}$  and has the composition  $Bi_{25}^{3+}Bi^{5+}O_{40}$ , where the two tetrahedral sites are occupied by one Bi3+ and one Bi5+ ion. Devalette *et al.* [10] reported that  $Bi_{12}[A_{1/2}^{3+}B_{1/2}^{5+}]O_{20}$ and  $Bi_{12}[Zn_{1/3}^{2+}B_{2/3}^{5+}]O_{20}$  (A = Fe, Ga and B = P, As, Bi) are derived from Bi<sub>12</sub>GeO<sub>20</sub> by couple substitutions. However, the existence of the Bi<sup>5+</sup> ion has not yet been verified directly. Accordingly, it is not clear at present whether the oxygen sublattice of the pure  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> is occupied completely or not.

The electrical conduction of bcc  $Bi_{24}Si_2O_{40}$  single crystal is reported to be p-type electronic at room temperature and to only have a small contribution from oxygen ionic transport at high temperatures [11, 12]. However,  $Bi_{24}Pb_2O_{38}$  and the nonstoichiometric  $Bi_{24}Si_2O_{40}$  sintered specimen were reported to have much higher oxygen ionic conductivities resulting in a mixed conduction behaviour [12, 13]. Harwig and Gerards [14] also suggested that the oxygen ionic conduction is predominant in the pure  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> based on the frequency dispersion of impedance. However, details of electrical conduction of the pure and doped  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> are not clear.

The present authors reported previously that the  $\gamma$ -form sintered specimen can be obtained at room temperature by doping a few mol % of Sb<sub>2</sub>O<sub>3</sub> [15, 16]. In the present study, the electrical conduction in  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> doped with Sb<sub>2</sub>O<sub>3</sub> was investigated by measuring electrical conductivity, as a function of temperature and oxygen partial pressure, and ionic transference number.

# 2. Experimental procedure

The raw materials were monoclinic  $\alpha$ -Bi<sub>2</sub>O<sub>3</sub> and Sb<sub>2</sub>O<sub>3</sub> (both 99.99% pure). The powders were mixed, calcined in air at 730° C for 2 h in a platinum crucible, finely ground, and pressed into pellets (10 mm in diameter  $\times$  3 mm in depth) or rods (5 mm  $\times$  3 mm  $\times$  10 mm) under 54 MPa. The pressed bodies were sintered at 800° C for 5 h in air and furnace-cooled.

Polymorphic forms and lattice constants were determined by X-ray diffraction (XRD). Polymorphic transformations were examined by high-temperature XRD and differential thermal analysis (DTA). The concentration of  $Sb^{3+}$  and  $Sb^{5+}$  ions in  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> were determined by wet chemical analysis.

Electrical conductivity was measured by a complex impedance method at 5 to 13 MHz and a four-probe d.c. method. The four-probe d.c. method was used for the determination of the oxygen partial pressure ( $P_{O_2}$ ) dependence of conductivity. No significant deviation in conductivity was observed for both methods. In order to control  $P_{O_2}$ , Ar–O<sub>2</sub> mixed gas was passed as a carrier gas and stabilized zirconia cells were used for reducing and monitoring  $P_{O_2}$ .

Ionic transference number was determined by a d.c.



Figure 1 Lattice constant of  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> doped with Sb<sub>2</sub>O<sub>3</sub> ( $\bullet$ ; [6]).

polarization method. A sputtered gold cathode was covered with indium metal as a blocking electrode, and a gold anode was exposed to air. The following equation was used to calculate the transference number  $t_i$ ;

$$t_i = 1 - \sigma(\infty)/\sigma(0) \tag{1}$$

where  $\sigma(0)$  and  $\sigma(\infty)$  are the d.c. conductivity at initiation of biassing and after polarization, respectively. The experimental set-ups are shown elsewhere [13, 17].

### 3. Results and discussion

Samples of b c c  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> were obtained by doping with 1 to 3 mol % Sb<sub>2</sub>O<sub>3</sub>. Minor monoclinic XRD peaks were detected only in the 1 mol % specimens. Doping with more than 3 mol % gave rise to another metastable tetragonal  $\beta$ -form. The apparent density of the sintered specimens in the  $\gamma$ -form was 88 to 92% of the theoretical value.

The lattice constant decreased with an increase of  $Sb_2O_3$  content as shown in Fig. 1. This indicates that  $Sb_2O_3$  dissolves into  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> in the present composition. Wet chemical analysis showed that more than 80% of antimony ions are pentavalent in 2 mol %  $Sb_2O_3$  doped  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub>. Fig. 2 shows DTA curves of  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> doped with 2 mol %  $Sb_2O_3$  and XRD patterns at the indicated temperatures. It was shown that  $\gamma$ -form transforms into  $\alpha$ -form at about 560°C and then into the  $\delta$ -form at

about 710° C on increasing temperature. On cooling, the  $\delta$ -form transformed directly into the  $\gamma$ -form.

Fig. 3 shows typical impedance plots of  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> doped with 32 mol % Sb<sub>2</sub>O<sub>2</sub> at 356 and 425° C. One semicircle in contact with the zero point at high frequencies and a section of semicircle at low frequencies can be observed. The low frequency region (below 1 kHz in Fig. 3) can be explained by electrode impedance [18]. There was no semicircle associated with the grain boundary, which often appears in middle frequency region for oxygen ionic conductors such as stabilized zirconia [19]. It is not clear whether this is due to the very small grain-boundary impedance or to the identical time constants in grain and grain-boundary impedances, which prevents the spliting of both contributions in the impedance plot. However, in the present study, the end-point of the first semicircle was taken as the bulk resistance to calculate conductivity.

The temperature dependences of the conductivity for  $Sb_2O_3$  doped  $Bi_2O_3$  on heating are shown in Fig. 4. The polymorphic transformations into the  $\alpha$  and  $\delta$ -forms are accompanied by abrupt conductivity changes. The composition dependence of the conductivity in the  $\gamma$ -form will be discussed later.

Ionic transference numbers in the Sb<sub>2</sub>O<sub>3</sub> doped Bi<sub>2</sub>O<sub>3</sub> are shown in Fig. 5. In the present method, the applied voltage was controlled to be below 50 mV. Accordingly, the measured transference numbers correspond to the average value under a  $P_{O_2}$  of between  $10^{4.33}$  (air) and  $10^{2.7}$  Pa. The transference numbers were almost unity in the b c c  $\gamma$ -form (below 550° C) and f c c  $\delta$ -form (above 720° C), whereas they were almost zero in the monoclinic  $\alpha$ -form. After the polarization measurements, morphological changes in the specimens were not found; those would be expected if charge transport by cations had occurred.

The dependence of conductivity on  $P_{O_2}$  for  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> doped with 1 to 3 mol % Sb<sub>2</sub>O<sub>3</sub> at 500° C are shown in Fig. 6. The conductivities were kept constant under the  $P_{O_2}$  of 10<sup>5</sup> to 10<sup>-9</sup> Pa. The results for the ionic transference number and  $P_{O_2}$  dependence of conductivity demonstrate that the  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> doped with 1 to 3 mol % Sb<sub>2</sub>O<sub>3</sub> is an oxygen ionic conductor in the wide range of  $P_{O_2}$ .



Figure 2 (a) DTA curve and (b) high-temperature XRD patterns at indicated temperatures for Bi<sub>2</sub>O<sub>3</sub> doped with 2 mol % Sb<sub>2</sub>O<sub>3</sub>.



Fig. 7 shows the dependence on Sb<sub>2</sub>O<sub>3</sub> content of conductivity  $\sigma$ , activation energy *E* and pre-exponential term  $\sigma_0$ , when they are expressed as  $\sigma = \sigma_0 \exp(-E/kT)$  where *k* is the Boltzmann constant and *T* the absolute temperature. The conductivity showed a maximum at 2.5 mol % Sb<sub>2</sub>O<sub>3</sub>. The activation energy remained almost unchanged and hence it is clear that the change of  $\sigma$  is due to that of  $\sigma_0$ .

It is proposed that pure  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> contains Bi<sub>26</sub>O<sub>39</sub> or  $Bi_{25}^{3+}Bi^{5+}O_{40}$  in the unit cell [1, 7–9]. However, if the oxygen sublattice is completely occupied,  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> would not show such a high conductivity. In fact, estimated oxygen ionic conductivity at 500° C in  $Bi_{24}Si_2O_{40}$  and non-stoichiometric  $Bi_{24}Si_{2-x}O_{40-2x}$ (x = 0.1 to 0.14) are reported to be  $3 \times 10^{-10}$  and  $4 \times 10^{-5} \Omega^{-1} \mathrm{cm}^{-1}$ , repectively [12, 13]. Those are much lower than that of  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> doped with Sb<sub>2</sub>O<sub>3</sub>. In addition, the reported density of pure  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub>, 9.21 g cm<sup>-3</sup> [7], is smaller than the calculated density of 9.29 for  $Bi_{26}O_{39}$  and 9.315 g cm<sup>-3</sup> for  $B_{26}O_{40}$  (using 1.0268 nm as the lattice constant), but larger than  $8.575 \,\mathrm{g}\,\mathrm{cm}^{-3}$  for  $\mathrm{Bi}_{24}\mathrm{O}_{36}$  assuming that all tetrahedrons are vacant. Accordingly, it is assumed that the oxygen sublattice in pure  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> is only partially occupied



Figure 4 Temperature dependence of conductivity of  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> doped with Sb<sub>2</sub>O<sub>3</sub>; ( $\bullet$ ) 1, ( $\triangle$ ) 1.5, ( $\blacktriangle$ ) 2, ( $\bigcirc$ ) 2.5, ( $\Box$ ) 3 mol % Sb<sub>2</sub>O<sub>3</sub>.

Figure 3 Complex impedance plots of  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> doped with 3 mol % Sb<sub>2</sub>O<sub>3</sub> at 356 and 425° C. Indicated numbers are the measuring frequency (Hz).

and consequently the large vacancy concentration brings about a high oxygen ionic conductivity. However, it is not clear whether those vacancies are caused only by  $Bi^{3+}$  occupation at tetrahedral  $Si^{4+}$  sites or by some Shottky-type defects of  $Bi^{3+}$  (or  $Bi^{5+}$ ) and oxygen vacancies [13].

As shown in Fig. 7, the pre-exponential term  $\sigma_0$  increased with increasing Sb<sub>2</sub>O<sub>3</sub> content up to 2.5 mol %. This indicates an increase in the oxygen vacancy concentration effective for ionic conduction [20]. Reported ionic radii (6-coordinated) are 0.103 and 0.076 nm for Bi<sup>3+</sup> and Bi<sup>5+</sup>, 0.076 and 0.060 nm for Sb<sup>3+</sup> and Sb<sup>5+</sup>, and 0.040 nm for Si<sup>4+</sup>, respectively, based on 0.140 nm for O<sup>2-</sup> [21]. The oxygen vacancies in  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> must be bound or ordered around Bi<sup>3+</sup> (or Bi<sup>5+</sup>) at tetrahedral sites, which are too large to occupy those sites. The incorporated smaller Sb<sup>5+</sup> ions would probably substitute preferentially at the tetrahedral sites and subsequently the concentration of the bound vacancies would be decreased resulting in an increase in  $\sigma_0$ .

On the other hand, the bcc lattice approaches the perfect structure as the oxidized Sb<sup>5+</sup> ions and corresponding oxygen ions are further incorporated. When incorporated Sb<sub>2</sub>O<sub>5</sub> content is 3.85 mol %, composition becomes stoichiometric  $Bi_{24}^{3+}(Bi^{3+}Sb^{5+})O_{40}$ , although its presence is not yet confirmed. Accordingly, the decreases in  $\sigma_0$  and  $\sigma$  at above 2.5 mol % Sb<sub>2</sub>O<sub>3</sub> shown in Fig. 7 appear to be due to a decrease in the total oxygen vacancy concentration.



*Figure 5* Ionic transference numbers in  $Sb_2O_3$  doped  $Bi_2O_3$ . Indicated numbers show  $Sb_2O_3$  content (mol %).



Figure 6  $P_{O_2}$  dependence of conductivity at 500° C for  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> doped with Sb<sub>2</sub>O<sub>3</sub>; (•) 1, (•) 2, (□) 3 mol % Sb<sub>2</sub>O<sub>3</sub>.

# 4. Conclusions

Body centred cubic  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub> doped with 1 to 3 mol % Sb<sub>2</sub>O<sub>3</sub> was stable up to 550° C but transformed into monoclinic  $\alpha$  and f c c  $\delta$ -Bi<sub>2</sub>O<sub>3</sub> at high temperatures.

In the Sb<sub>2</sub>O<sub>3</sub> doped  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub>, the oxygen ionic conduction was predominant in the  $P_{O_2}$  range of 10<sup>5</sup> to 10<sup>-9</sup> Pa. As the Sb<sub>2</sub>O<sub>3</sub> content increased, oxygen ionic conductivity increased up to 2.5 mol % Sb<sub>2</sub>O<sub>3</sub> reaching a maximum of  $1.8 \times 10^{-3} \Omega^{-1} \text{ cm}^{-1}$  at 500° C and



*Figure 7* Sb<sub>2</sub>O<sub>3</sub> content dependences of conductivity  $\sigma$  (at 400 and 500° C), activation energy *E* and pre-exponential term  $\sigma_0$ .

then decreased. However, the activation energy for ionic conduction remained almost unchanged.

It was confirmed that more than 80% of antimony ions are pentavalent in  $\gamma$ -Bi<sub>2</sub>O<sub>3</sub>. Those Sb<sup>5+</sup> ions appear to affect the concentraion of oxygen vacancies effective for ionic conduction.

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